Localized configurations in the AdS/CFT correspondence

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J. Jankowski Localized configurations in AdS/CFT

• Holographic principle

Quantum gravity in d + 1dimensions must have a number of d.o.f. which scales like that of QFT in d dimensions



't Hooft and Susskind '93

String Theory realization AdS/CFT correspondence
 Solutions of Einstein equations in d + 1 dimensions ↔
 States in strongly coupled QFT in d dimensions

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- Fascinating links between GR and QFT!
- Physically motivated interesting questions can be addressed by GR methods

Maldacena '97

• The Anti-de-Sitter spacetime can be given in a Poincare path coordinates

$$ds^2 = \frac{\eta_{\mu\nu} dx^\mu dx^\nu + dz^2}{z^2} \tag{1}$$

where z > 0 and boundary is at z = 0

• Metric (1) solves (d + 1)-dim. Einstein equations

$$R_{ab} - \frac{1}{2}g_{ab}R - \frac{d(d-1)}{2}g_{ab} = 0$$
 (2)

• This corresponds to the vacuum state in field theory

$$\langle T_{\mu\nu} \rangle = 0$$

 In this talk d = 4 i.e. AdS₄/CFT₃ correspondence → motivated by applications to solid state physics

Field - Operator correspondence

 $\bullet\,$ Field in the bulk $\longleftrightarrow\,$ Local operator on the boundary

 $\phi(x,z)\longleftrightarrow \mathcal{O}(x)$

- Solve Einstein+matter equations with an arbitrary Dirichlet boundary condition φ(x, z = 0) = φ₀(x)
- $\phi_0(x)$ sources operator $\mathcal{O}(x)$
- The near boundary (z
 ightarrow 0) expansion

$$\phi(x,z) \sim \phi_0(x) z^{d-\Delta} + \phi_1(x) z^{\Delta} + \dots$$
(3)

where $\dim[\mathcal{O}] = \Delta$

The expectation value is the subleading term

$$\phi_1(x) = \langle \mathcal{O}(x) \rangle \tag{4}$$

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Example: AdS₄ black hole

- Equilibrium state in field theory → black hole in the bulk Temperature in field theory → Hawking temperature
- Solve for metric \rightarrow read off energy density

$$g_{tt}(z) \sim -1 + z^3 g_{tt}^{(3)} + \dots$$
 (5)

with

$$\varepsilon_{\rm CFT} = \langle T_{tt} \rangle = \frac{3N^2}{8\pi} g_{tt}^{(3)} = \frac{16N^2\pi^3 T^3}{27}$$
 (6)

which is determined by conformal symmetry

• To compare free massless bosons

$$\varepsilon_{\rm free} = \frac{T^3 N^2 \zeta(3)}{\pi} \tag{7}$$

I. R. Klebanov, A. A. Tseytlin, Nucl. Phys. B 475, 164 (1996)

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Scalar field in AdS₄/CFT₃

- On top of that one can add matter fields to model physical systems
- Scalar field of mass $m^2=-2$ has the near-boundary expansion

$$\phi(x,z) \sim \phi_0(x)z + \phi_1(x)z^2 + \dots$$
 (8)

• Deformation of the theory by an operator $\mathcal{O}(x)$ of conformal dimension $\Delta = 2$, with the source term $\phi_0(x)$

$$\mathcal{L} = \mathcal{L}_{\rm CFT_3} + \int d^3 x \phi_0(x) \mathcal{O}(x)$$
(9)

and the expectation value $\phi_1(x) = \langle \mathcal{O}(x)
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Dirac δ -like source

• Consider local line-like source

$$\phi_0(x) = \eta \delta(x) \tag{10}$$

- Source φ₀(x) = cos(kx) was used to mimic a lattice
 G. T. Horowitz, J. E. Santos and D. Tong, JHEP 1207, 168 (2012)
- Try to obtain a lattice $\sum_n \delta(x na) \rightarrow$ holographic version of the Kronig-Penney model
- Physics of single defects \rightarrow both line-like and point-like

Goal:

Develop techniques to numerically solve Einstein's equations with boundary conditions (10)

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Discontinuous boundary conditions in AdS/CFT

- θ(x) and m = 0 → Janus solutions (analytical, 1D ODE)
 D. Bak, M. Gutperle and S. Hirano, JHEP 0305, 072 (2003)
- θ(x) and m = 0 at T > 0 → Janus black holes in d = 2 + 1 (numerical PDE, analytical)
 D. Bak, M. Gutperle and R. A. Janik, JHEP 1110, 056 (2011)
- $\delta(x)$ and $m^2 = -2$ with SUSY \rightarrow analytical and scale invariant E. D'Hoker, J. Estes, M. Gutperle and D. Krym, JHEP **0906**, 018 (2009)

Linearized solution

• Line defect linearized solution around empty AdS₄

$$\phi_0(x,z) = \frac{\eta z^2}{\pi (x^2 + z^2)}$$
(11)

- Easy numerical generalization to finite temperature/chemical potential
- For the $m^2 = -2$ case linearized solution preserves the 1+1 conformal symmetry along the defect, where the operator is sourced
- Define new variable $\tan \alpha = x/z$ then

$$\phi_0(x,z) = \frac{\eta}{\pi} \cos^2(\alpha) \tag{12}$$

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and use it for the backreacted solution at $\mathcal{T}=0$

- $\, \bullet \,$ Assume full solution depends only on α
- The metric will have an AdS slicing

$$ds^{2} = \frac{1}{A(\alpha)^{2}} \left(d\alpha^{2} + \frac{dr^{2} - dt^{2} + dy^{2}}{r^{2}} \right) \qquad \phi = \phi(\alpha) \quad (13)$$

- AdS boundaries at both sides of the defect will appear at $\alpha = \pm \alpha_0$ with $\alpha_0 \neq \pi/2$
- We found both numerical and perturbative analytical solutions parametrized by $\phi(lpha=0)$

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- $\partial_{\alpha}\phi(0) \neq 0 \rightarrow$ dynamically generated sources for the scalar!
- ullet Dynamically generated scale ightarrow conformal invariance breaking
- No reduction to an ODE system → logarithmic terms !?



Metric and scalar field for $\phi(0) = 0.9$ with N = 47 spectral points

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Regularized lattice

• Consider approximation of the Dirac δ lattice at T > 0 using linearized vacuum solution

$$\phi_P(x,z) = \sum_{n \in \mathbb{Z}} \phi_0(x+n,z) = \frac{\eta z \sinh(2\pi z)}{\cosh(2\pi z) - \cos(2\pi x)} \quad (14)$$

- Model of atomic lattice
- Very large gradients → numerically very difficult!



 $z_0 = 0.1, 0.15, 0.2$

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Grid modifications

 Analyze the structure of the linearized solution for the Dirac δ function



- Use this to adopt the grid points for numeric
- Analyze regularized solutions around the limiting case



- We observe that solutions depend on a scaling variable $w = \frac{z}{z_0}$
- In $z_0
 ightarrow 0$ limit only scale invariant derivatives appear $w \partial_w$
- Asymptotic $w \to \infty$ should carry information about the proper boundary conditions



We can observe evidence for logarithmic behavior of the backreacted ${\rm Dirac}\text{-}\delta$ lattice solution



- Example of a novel setup in Numerical Relativity
- Development of new numerical techniques to handle the problem → e.g. modified grids in spectral methods
- Detailed understanding of boundary conditions
- \bullet Holographic model of a lattice \rightarrow solid state physics e.g. optical conductivity, heat transport
- Work in progress ...

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