# Semiclassical limit and transfer matrix in Causal Dynamical Triangulation

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Andrzej Görlich Causal Dynamical Triangulation

- Foundations of Casual Dynamical Triangulation
- Emergence of background geometry
- Semiclassical limit
- Effective action
- Transfer matrix
- Onclusions

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# What is Causal Dynamical Triangulation?

Causal Dynamical Triangulation (CDT) is a background independent approach to quantum gravity.

$$\int \mathbf{D}[g] e^{iS^{EH}[g]} \quad \rightarrow \quad \sum_{\mathcal{T}} e^{-S^{R}[\mathcal{T}]}$$

CDT provides a lattice regularization of the formal gravitational path integral via a sum over causal triangulations.



# Path integral formulation of quantum gravity

- **Quantum mechanics** can be formulated in terms of *path integrals*: All possible trajectories contribute to the transition amplitude.
- General Relativity: gravity is encoded in space-time geometry.
- The role of a trajectory plays now the geometry of four-dimensional space-time.
- All space-time histories contribute to the transition amplitude.



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Sum over all two-dimensional surfaces joining the in- and out-state

# Causality

- **Causal D**ynamical Triangulations assume global proper-time foliation. Spatial slices (leaves) have fixed topology and are not allowed to split in time.
- Foliation distinguishes between time-like and spatial-like links.
- Such setup does not introduce causal singularities, which lead to creation of baby universes. (EDT)
- CDT defines the class of admissible space-time geometries which contribute to the transition amplitude.



- Discretization is a the standard regularization used in QFT.
- 4D simplicial manifold is obtained by gluing pairs of 4-simplices along their 3-faces.
- Spatial states are 3D geometries with a topology S<sup>3</sup>.
   Discretized states are build from equilateral tetrahedra.
- The metric is **flat** inside each 4-simplex.
- Length of time links *a<sub>t</sub>* and space links *a<sub>s</sub>* is constant.
- Curvature is localized at triangles.

#### Fundamental building blocks of Euclidean DT



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#### 3D spatial slices with topology $S^3$



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- The metric is **flat** inside each 4-simplex.
- Length of time links  $a_t$  and space links  $a_s$  is constant.
- Curvature is localized at triangles.

# Fundamental building blocks of 4D CDT

## Regge action

The Einstein-Hilbert action has a natural realization on piecewise linear geometries called Regge action

$$S^{E}[g] = -\frac{1}{G} \int \mathrm{d}t \int \mathrm{d}^{D}x \sqrt{g}(R-2\Lambda)$$

 $N_0$  number of vertices  $N_4$  number of simplices  $N_{14}$  number of simplices of type {1,4}  $K_0 \ K_4 \ \Delta$  bare coupling constants ( $G, \Lambda, a_t/a_s$ )



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• The partition function of quantum gravity is defined as a formal integral over all geometries weighted by the Einstein-Hilbert action.

$$Z = \int \mathrm{D}[g] e^{i S^{EH}[g]}$$

- To make sense of the gravitational path integral one uses the standard method of regularization - discretization.
- The path integral is written as a nonperturbative sum over all causal triangulations  $\mathcal{T}$ .
- Wick rotation is well defined due to global proper-time foliation.  $(a_t \rightarrow ia_t)$
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### Spatial slices

- The simplest observable giving information about the geometry, is the **spatial volume**  $n_i$  defined as a number of tetrahedra building a three-dimensional slice  $i = 1 \dots T$ .
- Restricting our considerations to the spatial volume *n<sub>i</sub>* we reduce the problem to one-dimensional quantum mechanics.





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### De Sitter space-time as background geometry

- In phase **C** the time translation symmetry is spontaneously broken and the distribution *n<sub>i</sub>* is bell-shaped.
- The average volume  $\langle n_i 
  angle$  is with high accuracy given by formula



• It describes Euclidean **de Sitter** space (*S*<sup>4</sup>), a classical **vacuum** solution.



## De Sitter space-time as background geometry

- In phase **C** the time translation symmetry is spontaneously broken and the distribution *n<sub>i</sub>* is bell-shaped.
- The average volume  $\langle n_i \rangle$  is with high accuracy given by formula

$$\langle n_i \rangle = H \cos^3\left(\frac{i}{W}\right)$$

• It describes Euclidean **de Sitter** space (*S*<sup>4</sup>), a classical **vacuum** solution.



#### Hausdorff dimension

The time coordinate *i* and spatial volume  $\langle n_i \rangle$  scale with total volume  $N_4$  as a genuine four-dimensional *Universe*,

$$t = N_4^{-1/4} i,$$
  
$$\bar{v}(t) = N_4^{-3/4} \langle n_i \rangle = \frac{3}{4\omega} \cos^3\left(\frac{t}{\omega}\right).$$



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# Spectral dimension

Simulations of the diffusion process allow to compute spectral dimension  $d_s$ .



Extrapolation of the results gives short and long range behavior

 $d_s(\sigma \to 0) = 1.95 \pm 0.10, \quad d_s(\sigma \to \infty) = 4.02 \pm 0.10,$ 

where  $\sigma$  is a fictitious diffusion time.

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## Minisuperspace model



- Classical trajectory  $\bar{v}(t)$  corresponds to Euclidean de Sitter space ( $S^4$ ), a spatially homogeneous and isotropic vacuum solution.
- $\bullet~$  We ,,freeze" all degrees of freedom except volume and assume that metric on  $S^3\times S^1$  spacetime has form

$$ds^2 = dt^2 + a^2(t) d\Omega_3^2, \quad v(t) = a^3(t)$$

• In this particular case, the Einstein-Hilbert action takes form

$$S = \frac{1}{G} \int \mathrm{d}t \int \mathrm{d}\Omega \sqrt{g}(R - 6\lambda)$$

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### Minisuperspace model



#### Quantum fluctuations

• We can measure the correlation matrix of spatial volume fluctuations around the classical solution  $\bar{n} = \langle n \rangle$ ,

$$\mathbf{C}_{ij} \equiv \langle (n_i - \langle n_i \rangle) (n_j - \langle n_j \rangle) \rangle$$

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$$S[n = \bar{n} + \eta] = S[\bar{n}] + \frac{1}{2} \sum \eta_i [\mathbf{C}^{-1}]_{ij} \eta_j + O(\eta^3), \quad [\mathbf{C}^{-1}]_{ij} = \frac{\partial^2 S[n]}{\partial n_i \partial n_i}$$





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## Effective action

- Minisuperspace action  $S[v] = \frac{1}{G} \int \frac{\dot{v}^2}{v} + v^{\frac{1}{3}} \lambda v dt$ • Discretization  $S[n] = \frac{1}{\Gamma} \sum_t \left( \frac{(n_{t+1} - n_t)^2}{n_{t+1} + n_t} + \mu n_t^{1/3} - \lambda n_t \right)$
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- The slicing can be refined by taking into account the {3,2} and {2,3} simplices which are located between layers t and t + 1 (simplices {4,1} and {1,4}).
- Half integer times t + 1/2 can be assign to  $\{3, 2\}$  simplices,

$$n_{t+\frac{1}{2}} \equiv \rho \cdot N_{32}(t), \quad n_t \equiv N_{41}(t).$$



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#### Volume profile for integer and half-integer t



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• Using the covariance matrix **C**, we can determine the effective action

$$S_{eff} = \sum_{t \in \mathbb{Z}} \left( + \frac{1}{\Gamma_1} \frac{\left(n_t - n_{t \pm \frac{1}{2}}\right)^2}{n_t + n_{t \pm \frac{1}{2}}} - \frac{1}{\Gamma_2} \frac{(n_t - n_{t+1})^2}{n_t + n_{t+1}} + V\left[n_t, n_{t+\frac{1}{2}}\right] \right)$$

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   N<sub>41</sub> repel each other. It is an entropic effect.
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$$\frac{1}{\Gamma} = \frac{1}{2\Gamma_1} - \frac{1}{\Gamma_2}.$$
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### Transfer matrix

The model of Causal Dynamical Triangulations is completely determined by transfer matrix  $\mathcal{M}$  labeled by 3D triangulations  $\tau$  of spatial slices.

$$Z = \sum_{\mathcal{T}} e^{-S^{\mathcal{R}}[\mathcal{T}]} \stackrel{!}{=} \operatorname{Tr} \mathcal{M}^{\mathcal{T}}$$
$$P^{(\mathcal{T})}(\tau_1, \dots, \tau_{\mathcal{T}}) = \frac{1}{Z} \langle \tau_1 | \mathcal{M} | \tau_2 \rangle \langle \tau_2 | \mathcal{M} | \tau_3 \rangle \dots \langle \tau_{\mathcal{T}} | \mathcal{M} | \tau_1 \rangle$$

Matrix element  $\overline{\langle \tau_1 \rangle} \mathcal{M}[\tau_2 \rangle$  denotes the transition amplitude from state  $|\tau_1\rangle$  to  $|\tau_2\rangle$ .



$$\mathcal{P}^{(T)}(\tau_1,\ldots,\tau_T) \equiv \frac{1}{Z} \langle \tau_1 | \mathcal{M} | \tau_2 \rangle \langle \tau_2 | \mathcal{M} | \tau_3 \rangle \ldots \langle \tau_T | \mathcal{M} | \tau_1 \rangle$$

• The effective action obtained from the covariance matrix

$$S[n] = \frac{1}{\Gamma} \sum_{t} \left( \frac{(n_{t+1} - n_t)^2}{n_{t+1} + n_t} + \mu n_t^{1/3} - \lambda n_t \right)$$

suggests existence of an effective transfer matrix labeled by the scale factor.

• Misleading to think of aggregated state  $|n\rangle = \sum_{\tau \sim n} |\tau\rangle$ • Density operator  $\rho(n) = |n\rangle\langle n| \equiv \sum_{\tau \sim n} |\tau\rangle\langle \tau|$ 

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$$\mathcal{P}^{(T)}(n_1,\ldots,n_T) \equiv \frac{1}{Z} \langle n_1 | \mathcal{M}[\underline{n_2}] \langle n_2 | \mathcal{M} | n_3 \rangle \ldots \langle n_T | \mathcal{M} | n_1 \rangle$$

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• Misleading to think of aggregated state

$$n \equiv \sum_{\tau \sim n} \tau$$

• Density operator  $\rho(n) = \overline{|n\rangle\langle n|} \equiv \sum_{\tau \sim n} |\tau\rangle\langle \tau|$ 

$$\mathcal{P}^{(T)}(n_1,\ldots,n_T) = \frac{1}{Z} \overline{\langle n_1 | M | n_2 \rangle} \langle n_2 | M | n_3 \rangle \ldots \langle n_T | M | n_1 \rangle$$

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Effective transfer matrix  $\langle n|M|m \rangle$  labeled by scale factor

$$\langle n|M|m \rangle = \mathcal{N}e^{-L_{eff}(n,m)}$$

$$S[n] = \sum_{t} L_{eff}(n_t, n_{t+1})$$

$$L_{eff}(n,m) = \frac{1}{\Gamma} \left[ \frac{(n-m)^2}{n+m} + \mu \left( \frac{n+m}{2} \right)^{1/3} - \lambda \frac{n+m}{2} \right]$$

Assuming that

$$P^{(T)}(n_1,\ldots,n_T) = \frac{1}{Z} \langle n_1 | M | n_2 \rangle \langle n_2 | M | n_3 \rangle \ldots \langle n_T | M | n_1 \rangle$$

we can measure elements of M

$$\langle n|M|m\rangle = \sqrt{P^{(2)}(n,m)} \text{ or } \langle n|M|m\rangle = \frac{P^{(3)}(n_1 = n, n_2 = m)}{\sqrt{P^{(4)}(n_1 = n, n_3 = m)}}$$

and check that it is consistent minisuperspace model

$$\langle n|M|m\rangle = \mathcal{N}e^{-\frac{1}{\Gamma}\left[\frac{(n-m)^2}{n+m} + \mu\left(\frac{n+m}{2}\right)^{1/3} - \lambda\frac{n+m}{2}\right]}$$



#### Kinetic part

The effective action consists of kinetic and potential terms,

$$\langle n|M|m\rangle = \mathcal{N}e^{-\frac{1}{\Gamma}\left[\frac{(n-m)^2}{n+m} + v\left(\frac{n+m}{2}\right)\right]}, \quad v(x) = \mu x^{1/3} - \lambda x + \delta x^{-\rho}$$

for n + m = c

The kinetic term causes a Gaussian behavior for n + m = c

$$\langle n|M|c-n\rangle = \mathcal{N}(c)e^{-\frac{(2n-c)^2}{k(c)}}, \quad k(c) = \Gamma \cdot c$$

 $\Gamma \approx 26.1$  is constant for all ranges of n.



#### Potential part

The effective action consists of kinetic and potential terms,

$$\langle n|M|m\rangle = \mathcal{N}e^{-\frac{1}{\Gamma}\left[\frac{(n-m)^2}{n+m}+\nu\left(\frac{n+m}{2}\right)\right]}, \quad \nu(x) = \mu x^{1/3} - \lambda x + \delta x^{-\rho}$$

minisuperspace possible curvature potential corrections

The potential part can be extracted from gathered data for n = m

$$\log\langle n|M|n\rangle = -\frac{1}{\Gamma} \left(\mu n^{1/3} - \lambda n + \delta x^{-\rho}\right) + \text{const}$$



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#### An effective model

Let's consider a toy model in which states are given by volume profiles  $\{n_t\}$  rather than triangulations,

$$P^{(T)}(n_1,\ldots,n_T) \propto \langle n_1|M|n_2 \rangle \langle n_2|M|n_3 \rangle \cdots \langle n_T|M|n_1 \rangle,$$

where the transfer matrix elements are given by

$$\langle n|M|m\rangle = \mathcal{N}e^{-\frac{1}{\Gamma}\left[rac{(n-m)^2}{n+m}+v\left(rac{n+m}{2}
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ight]}, \quad v(x) = \mu x^{1/3} - \lambda x + \delta x^{-\rho}.$$

It reproduces the full CDT results very accurately, supporting the transfer matrix approach.



# Summary

- The model of Causal Dynamical Triangulation is manifestly diffeomorphism-invariant. No coordinates are introduced and only geometric invariants, like length or angle, are involved.
- Phase diagram consists of three phases. In phase C emerges a four-dimensional universe with well defined time and space extent.
- The background geometry corresponds to the Euclidean de Sitter space, i.e. classical solution of the minisuperspace model.
- The effective action obtained from both the correlation matrix and transfer matrix corresponds to the minisuperspace action and properly desribes the emergent background geometry and quantum fluctuations. In CDT no degrees of freedom are frozen.
- The semiclassical limit is a truly nonperturbative limit. The contribution coming from the action is as important as the entropy of geometric configurations.

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## Thank You!

Based on:

J. Ambjørn, J. Gizbert-Studnicki, A. Görlich, J. Jurkiewicz, T. Trześniewski, Nucl. Phys. B 849, 144 (2011).

J. Ambjørn, J. Gizbert-Studnicki, A. Görlich, J. Jurkiewicz, JHEP 09 (2012) 017

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